

OBJECTIVES

Description of the structure of nuclei as self-bound quantum many-body systems using realistic 2- and 3-body forces between protons and neutrons

$$\hat{\mathbf{H}} = \sum_{i < j} \frac{(\vec{p_i} - \vec{p_j})^2}{2 m A} + \sum_{i < j} V_{ij} + \sum_{i < j < k} V_{ijk} + \dots$$

- Ground state (gs) and excitation energies
- Radii and electromagnetic moments

NUCLEAR INTERACTION

Chiral Effective Field Theory (χ EFT) up to next-to-next-to-leading order (N²LO)

	Two-nucleon force	Three-nucleon force			
LO (Qº)					
NLO (Q²)					
N²LO (Q³)					

- Low-energy constants (LECs) for 2-body potential fitted to NN scattering data [1]
- Two additional LECs for 3-body force fitted to A = 3 system [2]

Phenomenogical 2-body potential: **Daejon16**

- Based on chiral EFT potential at N³LO
- Off-shell behavior fitted to 10 energy levels from A = 4 to A = 16 [3]

NO-CORE CI

Solve eigenvalue problem for wave function Ψ

 $\hat{\mathbf{H}} \Psi(r_1, \dots, r_A) = \lambda \Psi(r_1, \dots, r_A)$

by expanding the *A*-body wave function Ψ in Slater Determinants of single-particle HO states ϕ

$\Psi(r_1,\ldots,r_A)$	=	$\sum c_i \Phi_i$	(r_1,\ldots,r_A))	
$\Phi_i(r_1,\ldots,r_A)$	—	$\frac{1}{\sqrt{(A!)}}$	$\phi_{i1}(r_1) \ dots \ dots \ \phi_{i1}(r_A) \ \phi_{i1}(r_A)$	• • •	$\left. \begin{array}{c} \phi_{iA}(r_1) \\ \vdots \\ \phi_{iA}(r_A) \end{array} \right $

- Truncate on total number of HO excitations, such that Center-of-Mass motion factorizes
- Calculate physical observables $\langle \Psi_f | \hat{\mathcal{O}} | \Psi_i \rangle$
- Increase basis until observables converge

ABINITIO NUCLEAR STRUCTURE CALCULATIONS OF ATOMIC NUCLEI UP TO OXYGEN-16

Results for Energy Levels [4]



- Energies extrapolated to complete basis (see poster by J.P. Vary using ANN for extrapolation)
- χ EFT: gs energies within chiral truncation error estimate (grey bands) through A = 12, but ¹⁶O is significantly overbound [2]
- **Daejon16**: good agreement for gs energies through A = 14, but ¹⁶O slightly overbound
- Excitation energies: difference of extrapolated energies
- Natural parity spectra in agreement with experiment up to at least A = 12
- Unnatural parity states too high with χEFT , but well-described by **Daejon16**

MAGNETIC MOMENTS [1]

Results with LO, NLO, and N²LO potentials using LO M1 operator, with chiral truncation estimates



- Without χ EFT corrections to M1 operator, nearly converged in chiral expansion
- No such corrections up to N²LO for isospin T=0 states ⁶Li and ¹⁰B, which agree with data
- Discrepancy with data for T > 0 states consistent with estimates of meson-exchange contributions to M1 operator at N²LO

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CONCLUSIONS

- Good description of more than 70 states, excluding isobaric analog states
- Agreement with data improves noticably when select many-body input is used in determining the nuclear interaction
- χEFT potential fitted to A = 2 and 3 only agreement with data decreases for $A \ge 12$
- **Daejon16** fitted to 10 energy levels in *p*-shell gives better agreement above A = 12and for unnatural parity states

WORK IN PROGRESS

- Improved, higher-order χ EFT interactions
- Consistent electroweak operators
- Electroweak transitions, including (neutrinoless) double- β decay
- Heavier nuclei using effective Shell Model interactions based on NCCI calculations for A = 17 and 18 nuclei

Eigenvalue problem for large sparse matrix. Convergence of energies, radii, moments, decays, etc, require dimensions well over 10^{10} .

ITERATIVE SOLVER [5]

Locally Optimized Block Preconditioned Conjugate Gradient (LOBPCG): SpMV acting on block of vectors, which improves cache performance, allows for vectorization, and needs significantly less iterations compared to Lanczos algorithm



[1]	S.
[2]	E.
[3]	A. P. 1
[4]	P. 2 P. 2
[5]	M. an
[6]	H. Di

COMPUTATIONAL CHALLENGES





- SpMV most timeconsuming kernel
- Memory bound
- Use only half of symmetric matrix for both SpMV and $SpMV^{T}$

DISTRIBUTED SPMV/SPMM [6]

• Compressed Sparse Row works okay for SpMV (left), but inefficient for SpMV^T (right) in hybrid MPI/OpenMP applications • Compressed Sparse Block improves data locality, and allows for efficient OpenMP parallelization of both SpMV and SpMV^T



Binder *et al*, Phys. Rev. C98, 014002 (2018) Epelbaum *et al*, Phys. Rev. C99, 024313, (2019)

.M. Shirokov, I.J. Shin, Y. Kim, M. Sosonkina, Maris, and J.P. Vary, Phys. Lett. B761, 87 (2016)

Maris, arXiv:1906.03703 [nucl-th];

Maris, I.J. Shin, and J.P. Vary, in preparation (2019) . Shao, H.M Aktulga, C. Yang, E.G. Ng, P. Maris, nd J.P. Vary, Comp. Phys. Comm. 222, 1 (2018)

.M. Aktulga *et al*, IEEE Transactions on Parallel and istributed Systems 28, 1550 (2016)