

# A Sinc DVR Grid-Orbital Hybrid Basis for Electronic Structure



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## INTRODUCTION

- The main problem we wish to solve is to diagonalize a general molecular Hamiltonian with the non-relativistic Born-Oppenheimer approximation

$$\hat{H} = \hat{T} + \hat{V} = -\frac{1}{2} \sum_{i=1}^{N_e} \nabla_i^2 - \sum_{i=1}^{N_e} \sum_{k=1}^{N_{nuc}} \frac{Z_k}{|\mathbf{r}_i - \mathbf{R}_k|} + \sum_{i < j \leq N_e} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}$$

- Grid-orbital hybrid basis developed for single and double ionization applications
- Orbitals used to represent bound electrons and grid functions used to represent continuum electrons

## PRIMITIVE GRID BASIS

- The 1D grid is uniformly spaced with  $2n + 1$  points

$$x_j = j\Delta, \quad j = -n, \dots, n, \quad x_j \in [-n\Delta, n\Delta]$$

- Expand the wave function in the  $(2n + 1)$ -dimensional basis

$$\Psi(x) \approx \sum_{k=-n}^n \psi_k \tilde{\chi}_k(x),$$

$$\tilde{\chi}_k(x) = \frac{1}{\sqrt{\Delta}} \text{sinc}\left(\frac{x - x_k}{\Delta}\right), \quad \text{sinc}(x) = \frac{\sin(\pi x)}{\pi x}$$

- $L^2$  orthonormal:  $\langle \tilde{\chi}_i | \tilde{\chi}_j \rangle = \delta_{ij}$
- DVR Property:  $\tilde{\chi}_i(x_j) = \Delta^{-1/2} \delta_{ij}$
- Simple kinetic energy matrix elements known exactly

$$\langle \tilde{\chi}_i | \hat{T} | \tilde{\chi}_j \rangle = T_{ij} = \begin{cases} \pi^2 / (6\Delta^2) & \text{if } i = j \\ (-1)^{i-j} / (\Delta^2 (i-j)^2) & \text{if } i \neq j \end{cases}$$

- Potential energy matrix representation is diagonal w.r.t. DVR quadrature

$$\langle \tilde{\chi}_i | \hat{V} | \tilde{\chi}_j \rangle \approx \sum_{k=-n}^n \Delta \tilde{\chi}_i(x_k) V(x_k) \tilde{\chi}_j(x_k) = \delta_{ij} V(x_i)$$

- 3D grid functions defined as products of 1D functions

$$\chi_{ijk}(\vec{r}) = \chi_i(x) \chi_j(y) \chi_k(z)$$

## PRIMITIVE HYBRID BASIS

- Let  $M_1 = (2n + 1)^3 =$  dimension of the grid basis
- The one-electron hybrid basis consists of a set of  $M_2 \ll M_1$  orthonormal orbitals expanded in the grid basis

$$\phi_\alpha(x) = \sum_{i=1}^{M_1} C_{i\alpha} \tilde{\chi}_i(x), \quad \alpha = 1, \dots, M_2$$

and a set of modified grid functions defined as

$$\chi_i(x) = \sum_{j=1}^{M_1} Q_{ij} \tilde{\chi}_j(x), \quad i = 1, \dots, M_1$$

- The matrix  $Q = 1 - CC^T$  is chosen to ensure that every modified grid function is orthogonal to every orbital
- The hybrid basis  $\mathcal{B}$  is defined by combining the modified grid functions and the orbitals

$$\mathcal{B} = \{\xi_1, \dots, \xi_M\} = \{\chi_1, \dots, \chi_{M_1}, \phi_1, \dots, \phi_{M_2}\}, \quad M = M_1 + M_2$$

- Orbitals generated with Hartree-Fock method

## HYBRID BASIS MATRIX ELEMENTS

- Overlap matrix elements:

$$\langle \chi_i | \chi_j \rangle = Q_{ij}, \quad \langle \phi_\alpha | \phi_\beta \rangle = \delta_{\alpha\beta}, \quad \langle \chi_i | \phi_\alpha \rangle = 0$$

- Denote the one and two electron matrix elements of the primitive grid basis as

$$\tilde{h}_{ij} = \langle \tilde{\chi}_i | \hat{h} | \tilde{\chi}_j \rangle, \quad [\tilde{j} | \tilde{k} \tilde{\ell}] = \delta_{ij} \delta_{kl} \tilde{V}_{ik}$$

- One-electron Hamiltonian matrix elements:

$$\langle \chi_i | \hat{h} | \chi_j \rangle = \sum_{k,\ell} Q_{ik} Q_{j\ell} \langle \tilde{\chi}_k | \hat{h} | \tilde{\chi}_\ell \rangle = \sum_{k,\ell} Q_{ik} Q_{j\ell} \tilde{h}_{k\ell}$$

$$\langle \chi_i | \hat{h} | \phi_\alpha \rangle = \sum_{k,\ell} Q_{ik} C_{\ell\alpha} \langle \tilde{\chi}_k | \hat{h} | \tilde{\chi}_\ell \rangle = \sum_{k,\ell} Q_{ik} C_{\ell\alpha} \tilde{h}_{k\ell}$$

$$\langle \phi_\alpha | \hat{h} | \phi_\beta \rangle = \sum_{k,\ell} C_{k\alpha} C_{\ell\beta} \langle \tilde{\chi}_k | \hat{h} | \tilde{\chi}_\ell \rangle = \sum_{k,\ell} C_{k\alpha} C_{\ell\beta} \tilde{h}_{k\ell}$$

- Two-electron Hamiltonian matrix elements:

$$[ij]k\ell] = \sum_{i'k'} Q_{i'i'} Q_{j'j'} Q_{k'k'} Q_{\ell'\ell'} \tilde{V}_{i'k'}$$

$$[i\alpha]j\beta] = \sum_{i'k'} Q_{i'i'} C_{i'\alpha} Q_{j'j'} C_{k'\beta} \tilde{V}_{i'k'}$$

$$[ij]k\alpha] = \sum_{i'k'} Q_{i'i'} Q_{j'j'} Q_{k'k'} C_{k'\alpha} \tilde{V}_{i'k'}$$

$$[i\alpha]j\beta\gamma] = \sum_{i'k'} Q_{i'i'} C_{i'\alpha} C_{k'\beta} C_{k'\gamma} \tilde{V}_{i'k'}$$

$$[ij]k\alpha\beta] = \sum_{i'k'} Q_{i'i'} Q_{j'j'} C_{k'\alpha} C_{k'\beta} \tilde{V}_{i'k'}$$

$$[\alpha\beta]j\gamma\delta] = \sum_{i'k'} C_{i'\alpha} C_{i'\beta} C_{k'\gamma} C_{k'\delta} \tilde{V}_{i'k'}$$

## MULTI-ELECTRON BASIS

- The  $N$ -electron wave function is expanded in a basis of Slater determinants

$$\Psi(x_1, \dots, x_N) = \sum_{\vec{m}} c_{\vec{m}} |\vec{m}\rangle$$

- The  $M$  hybrid basis functions are used as the orbitals in the determinants

$$|\vec{m}\rangle = |\xi_{n_1} \dots \xi_{n_N}\rangle, \quad 1 \leq n_i \leq M$$

- We only allow configurations with either 0,1 or 2 modified DVRs

$$|ij\vec{m}\rangle \quad (2 \text{ modified DVRs, } N - 2 \text{ orbitals})$$

$$|i\vec{m}\rangle \quad (1 \text{ modified DVR, } N - 1 \text{ orbitals})$$

$$|\vec{m}\rangle \quad (0 \text{ modified DVRs, } N \text{ orbitals})$$

- The  $\vec{n}$  configurations are generated by taking single and double excitations of the Hartree-Fock configuration
- For single ionization applications, we do not include the determinants with two grid functions
- The overlap matrix has a block diagonal structure

$$\mathcal{S} = \begin{bmatrix} \langle ij\vec{n} | i'j'\vec{n}' \rangle & 0 & 0 \\ 0 & \langle i\vec{n} | i'\vec{n}' \rangle & 0 \\ 0 & 0 & \langle \vec{n} | \vec{n}' \rangle \end{bmatrix}, \quad \begin{aligned} \langle ij\vec{n} | i'j'\vec{n}' \rangle &= \delta_{\vec{n}\vec{n}'} Q_{i'i'} Q_{j'j'} \\ \langle i\vec{n} | i'\vec{n}' \rangle &= \delta_{\vec{n}\vec{n}'} Q_{i'i'} \\ \langle \vec{n} | \vec{n}' \rangle &= \delta_{\vec{n}\vec{n}'} \end{aligned}$$

- Since the basis functions are non-orthogonal, the typical Slater-Condon rules for Hamiltonian matrix elements do not apply and we develop a customized set of rules for determining these elements
- The method is made practical due to the low rank of the projector  $P = CC^T$  and the fact that all orbitals are orthogonal to all grid functions

## HELIUM GROUND STATE

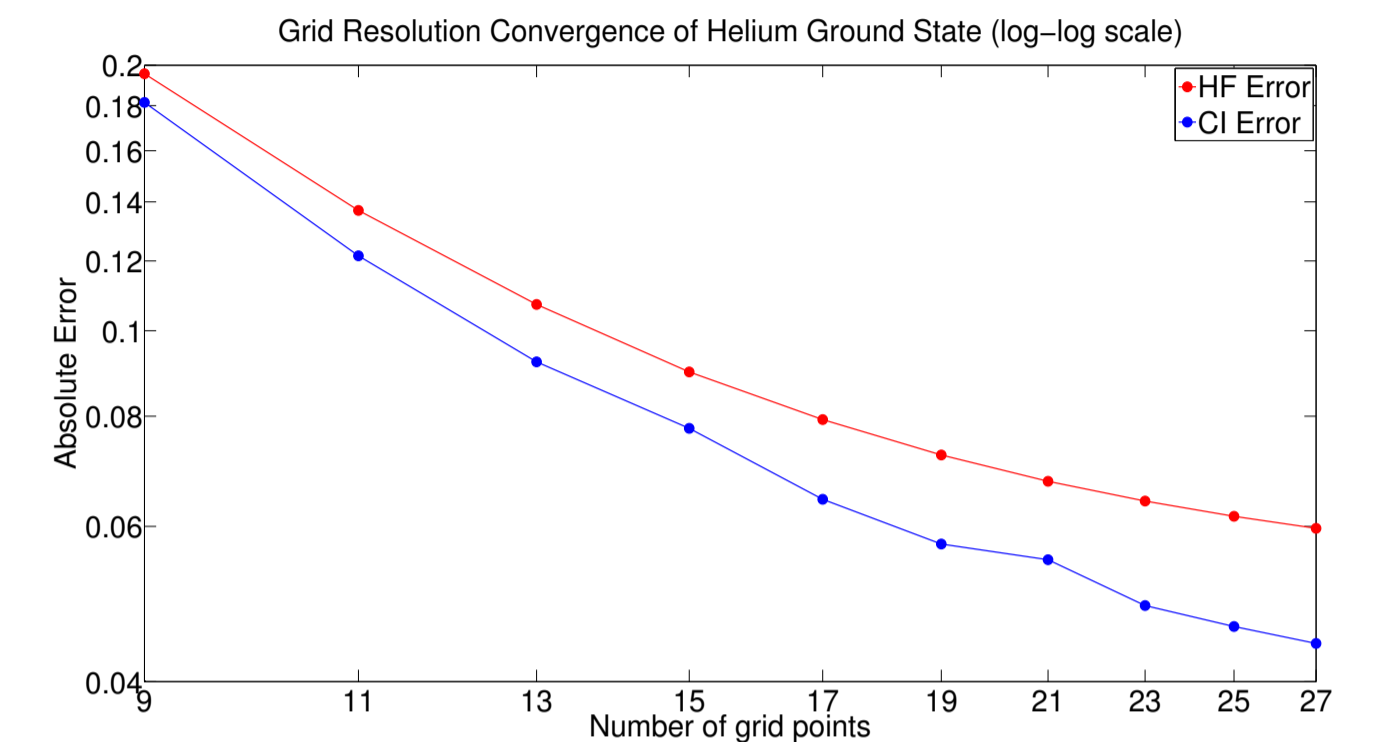


Figure: The absolute error in the calculated helium ground state energy for the Hartree-Fock method and the CI hybrid basis with one electron on the grid. The calculation used 8 spin orbitals and a fixed box size of 3 bohr.

## H<sub>2</sub> GROUND STATE

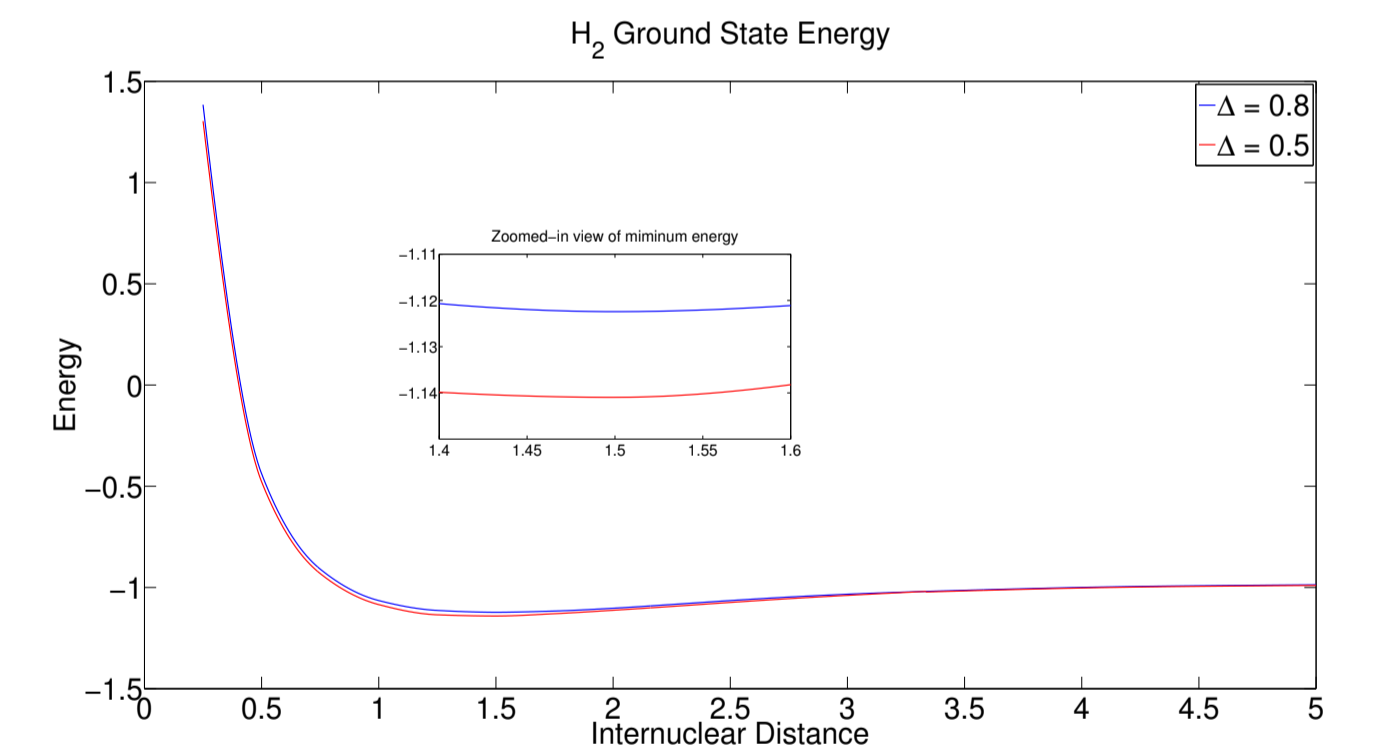


Figure: Ground state energy of H<sub>2</sub> vs. internuclear distance with one electron on the grid for  $\Delta = 0.5, 0.8$  and using 8 spin orbitals.

## WORK IN PROGRESS

- Implement MCSCF method in the sinc DVR basis to obtain better orbitals
- Include 2 electrons on the grid for double-ionization applications
- Incorporate exterior complex scaling
- Implement a multi-resolution method, different grids for different orbitals