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Lattice QCD on the BGQ: Achieving 1 PFlops Production Jobs

Theory of

interactions

of quarks

Known Elementary Particles QCD

Interactions

mediated

by gluons

 \bar{d}

c

QCD + Electroweak Interactions Produce Particle Decays \overline{O} mixings of \overline{O} mixings \overline{O} . The Standard Model (SM). \sqrt{U} arise from the \sqrt{U}

 $\frac{\frac{2}{3}}{\frac{1}{2}}$ U $\frac{2}{3}$ C $\left\| \cdot \right\|$

 \overline{u}

d

 \overline{u}

s

s

s

• Use multimass Krylov space solver (CG) to do all poles at once

as as $f+(0)$ s12s23−c12s23−c12s23−c12s23−c12s23−c12s23−c12s23−c12s23−c12s23−c12s23−c12s23−c12s23−c12s23−c12s23−c12s23−c12s23
s13ei5 c23c13 s13ei5 c23c

For K_{l3} we have: $\Gamma_{K \to \pi l \nu} = C_K^2 \frac{G_F^2 m_K^3}{192 \pi^3} I S_{EW} [1 + 2 \Delta_{SU(2)} + 2 \Delta_{EM}] |V_{us}|^2 |f_+(0)$ 2 3 $\frac{2}{5}$ m 2 $2 |f(\Omega)|^2$ π For K_{3} we have: $\Gamma_{K\to \pi l\nu} = C_{K}^{2} \frac{G \bar{F} mK}{100 \pi^{3}} I S_{EW} [1 + 2 \Delta_{SU(2)} + 2 \Delta_{EM}] |V_{us}|^{2} |f_{+}$

- QCD Simulations are done in a four (or five) dimensional box, with O(50-100) grid points in each dimension Ω OCD Simulations are done in this hierarchy using the Wolfenstein parameterization. We define QCD Simul n each Ĭ
	- Sample the phase space of the system via Monte Carlo, following the Euclidean space Feynamn path integral. The continuum, Minkowski space path integral is $\frac{1}{2}$. Sample $\frac{1}{2}$ the pl \mathbf{I} \mathbf{R}

 $Z =$ \overline{dA} 3 $i=1$ $Z = \int [dA] \prod_{i=1}^{3} \det [D(A, g_0, m_0^i)] \exp \left\{-\frac{i}{4}(\partial_\mu A_\nu - \partial_\nu A_\mu + g_0 f^{abc} A^b_\mu A^c_\nu)^2\right\}$ $D(A, g_0, m_0^i) \equiv i \gamma^\mu (\partial_\mu - i g A_\mu^a t^a/2) - m_0^i$ $Z = \int [dA] \prod \det [D(A)]$ \mathcal{J} matrix $\sum_{i=1}^{\bullet}$ $\mathcal{D}(A, \alpha, m^{i}) = i\alpha l^{\mu}$ $L(1, g_0, m_0) = v_0$

> • Only parameters are an input coupling and three (or four) light quark masses. $\overline{}$ $\overline{\$

> > Numerical Path Integral Including the Fermionic Effects of Quarks

• Require positive definite measure for Monte Carlo

 $Z =$ $\int [dA]$ \prod $\scriptstyle i=u,d,s$ $\left(\frac{\det[D^{\dagger}(m_i)D(m_i)]}{\det[D^{\dagger}(1)D(1)]}\right)^{1/2}\exp\left\{-S_g(A)\right\}$ = \mathcal{L} $[dA][d\Phi^\dagger_i] [d\Phi_i]$ exp $\sqrt{ }$ $\bigg\{-S_g(A)-\sum_{i=u,d,s}$ $\Phi_i^{\dagger}[D(1)\,D^{\dagger}(m_i)^{-1}\,D(m_i)^{-1}\,D^{\dagger}(1)]^{1/2}\Phi_i$ \mathbf{A} \mathbf{I} $\frac{1}{2}$

RBC/UKQCD have production jobs on the Argonne ALCF BGQ that sustain 1 PFlops on 32 racks = $32k$ nodes = 0.5 M cores.

Algorithms for Measurements $\mathbf{A} \mathbf{1} \quad \mathbf{A} \mathbf{1} \quad \mathbf{A} \quad$

This performance comes from very carefully tuned assembly code on BGQ, produced by Peter Boyle, using his BAGEL code generator

Algorithms for Gauge Field Production

- Producing gauge fields:
- * Use classical molecular dynamics to move through gauge field space
- * Quark loops give back reaction on gauge fields by solving Dirac equation
- asenbusch mass p

- sufficiently decorrelated (independent enough) to make them worth calculating <u>1 (</u> • Time translated the n-point function, on a fixed background gauge field, are
- This means many solutions of the Dirac equation $D[U_\mu] \Psi = s$ for fixed U_μ
- Calculating eigenvectors of $D[U_{\mu}]$ with small eigenvalues (low-modes) speeds up subsequent solves. Can be done with EigCG or Lanczos algorithms with twisted kaon. Bottom left: operator syn \mathbb{R}^n with twisted pion. Bottom right: operator \mathbb{R}^n
- Alternatives for Wilson fermions are domain decomposition and multigrid, giving similar speed-up with smaller memory requirements.
- Further improvement from all-mode-averaging of Blum, Izubuchi and Shintani
- * Separates measurements into expensive parts, with small statistical errors after a

0 10 20 30 40 50 60

t

 $T = 30$ $T_{\rm i} = 40$ $T = 50$ $T = 60$

0 10 20 30 40 50 60

- few measurements, and inexpensive parts, where many measurements are needed.
- These improvements make measurements $\sim 10 \times$ faster than a year ago.
- * This particular method takes substantial computer memory O(100 TBytes)
- * No checkpointing is done, so computer must work reliabilty for duration of job
- * Our smaller volume simulation takes 6 days on 1 rack of BGQ
- * Our larger volume simulations take 6 hours on 32 BGQ racks at 1 PFlops.
- * This speed-up was accomplished through the work of Hantao Yin at Columbia.
- Boyle has developed a Heirarchically Deflated Conjugate Gradient (HDCG) to markedly decrease the memory footprint, since eigenvectors are not needed, while also needing fewer iterations to converge. Production level testing is currently underway by Chulwoo Jung of BNL

 t the simulations, the lattice spacing and the physical values \mathcal{L} of the quark masses can only be estimated. Their precise

Through advances in computer hardware and software, production lattice QCD jobs run by the RBC and UKQCD Collaborations on BGQ installations are sustaining 1 PFlops. The largest computational cost is in the solution of the Dirac equation (a linear equation) in the presence of a fixed gauge background and we use Peter Boyle's (Univ. of Edinburgh) BAGEL based solver for this important calculation. Extensive multithreading via OpenMP in the rest of our code base (Chulwoo Jung, BNL and Hantao Yin, Columbia) has also been vital to achieving good performance. In addition, we employ a number of new deflation and variance reduction strategies in our calculations which give a further large speed-up $($ \sim 10 \times) beyond the software improvements. With all of these techniques, we are able to simulate lattice QCD with physical pion masses in large volumes, of size $(5.5 \text{ fm})^3$. This has led to markedly reduced statistical and systematic errors for our results. With these advances, as well as new theoretical ideas, we are now begining production calculations for kaon decay involving disconnected quark diagrams, where the signal to noise ratio is much worse. These new calculations improve constraints on standard model physics.
Sequoia at LLNL
The coloniations reported

- \bullet Decays of quarks via weak interactions predicted by Standard Model.
- Experiments measure decays of hadrons
- Standard Model quark decays involve • Standard Model quark decays involve
elements of a 3 by 3 unitary matrix, the CKM matrix, described by 4 parameters m_A terms for the physical states are obtained by diagonalizing α u,d α ^L ^Y ^f ^V ^f†

← Gluon self-interactions yields a very non-linear system. Aλ3(1 [−] ^ρ [−] iη) [−]Aλ² ¹

 \overline{I} $\overline{10}$ physicists are now requilared • Given advances in hardware, software and algorithms, lattice QCD physicists are now regularly costhe beginning of numerical doing simulations with physically light pions, a major goal since the beginning of numerical QCD in the early 1980's.

paper. Note that we have generated three ensembles with a 0:12 fm and ml ¼ ms=10 that differ only in their

• Simulations are just beginning for the more difficult case of $K \rightarrow \pi \pi$ decays with the pions in an isospin 0 final state. This process includes disconnected diagrams, which are much noisier

- HDCG or other linear solver improvements could be very helpful for these measurements
- Simulations with a smaller underlying lattice spacing pose problems with the rate of sampling of the path integral phase space. Greg McGlynn at Columbia has made some progress on this important question.

$\overline{\mathcal{C}}$ * Hasenbusch mass preconditioning allows tuning back reaction

fermion actions, each lattice fermion species corresponds to four ''tastes'' of fermions in the continuum limit. To eliminate the three unwanted tastes from the quark sea, we use the fourth-root procedure for each of the sea-quark flavors, up, down, strange, and charm. For numerical and theoretical arguments justifying this fourth-root procedure,

- ses 7 levels of inter ICI Z^V * RBC/UKQCD uses 7 levels of intermediate masses
- * Integrate different d.o.f on different time scales (Sexton-Weingarten integrators)
- \mathbf{r} generate both the initial and final pions. The setup is shown in the right panel of figure * Use higher order integrators, currently RBC/UKQCD use force gradient, $O(dt^4)$
- These are giving 10-100× speed-up over a decade ago.

 $T = |t_K - t_\pi|$

* Hard to be completely quantitative here, since without speeds-ups, we could not even try current simulations computed at α

• Previous results

 $\mathrm{Re}\,A_2~=~(1.381\pm 0.046_{\textrm{\tiny stat}}\pm 0.135_{\textrm{\tiny sys}\,\textrm{\tiny no}\,a^2}\pm 0.207_{\textrm{\tiny a}^2})~\,\times 10^{-8}~~\mathrm{GeV}$ ${\rm Im} \ A_{\scriptscriptstyle 2} \ = \ \ - (6.54 \pm 0.46_{\scriptscriptstyle\rm stat} \pm 0.72_{\scriptscriptstyle\rm sys\,no\,a^2} \pm 0.98_{\scriptscriptstyle a^2}) \quad \times 10^{-13} \ \ {\rm GeV}$ 8 13 $\sigma = (1.381 \pm 0.046_{\rm stat} \pm 0.135_{\rm sys\, no\, a^2} \pm 0.207_{\rm a^2})\ \ \times 10^{-3}$ $\hspace{1.6cm} = \hspace{.4cm} - (6.54 \pm 0.46_{\scriptscriptstyle\rm stat} \pm 0.72_{\scriptscriptstyle\rm sys\,no\,a^2} \pm 0.98_{\scriptscriptstyle\rm a^2}) \hspace{.3cm} \times 10^{-7}$

• New, preliminary results

 $.424 \pm 0.041_{stat} \pm 0.135_{sys\, no\, a^2} \pm 0.13$ $.88 \pm 0.16$ _{stat} ± 0.72 _{sys no $a^2 \pm 0$.} Re Im *A A* $1.424 \pm 0.041_{\text{stat}} \pm 0.135_{\text{sys no }a^2} \pm 0.0$ 5.88 ± 0.16 stat \pm 0.72 sys no a^2 \pm 0.0 10 10 **GeV** $G_{\alpha^2} \pm 0.0_{\alpha^2}$ $\times 10^{13}$ GeV (Preliminary) 2 2 8 13 stat \perp \cup . 1 \cup \cup sys no stat \perp **U.** \angle sys no 2 \pm U.U.² 2 \pm U.U.² \pm 0.041_{stat} \pm $0.135_{\text{sys no }a^2}$ \pm \pm 0.16 stat \pm 0.72 sys no $a^2 \pm$ \times \times = $=$ $-$ - - $\big($ $\big($ $\big)$ $A_2 = -(5.88 \pm 0.16_{\text{stat}} \pm 0.72_{\text{sys no }a^2} \pm 0.0_{a^2})$

• These new measurements are reducing the errors on measurements of properties of decay of kaons into a pion plus leptons, called Kl3 decays **(RBC+UKQCD Collaborations)** $T_{\rm F}$ reprofits, called type accupations, $T_{\rm F}$

 $\langle \pi | V | K \rangle \rightarrow f_{+}^{K \pi} (q^2 = 0)$ • part. twisted boundary conditions • $N_f = 2 + 1$ domain wall fermions \bullet a^2 -scaling study (0.09fm-0.14fm) \rightarrow tiny cut-off effects • physical point simulation *^m*π: 171–670MeV [→] arXiv:1305.7217 137–670MeV → PRELIMINARY • polynomial ansatz describes data over entire mass range • phys. point data eliminates large systematic due to χ extrapolation

$f_+^{K\pi}(0)=0.9670\big(20\big)\!\big(\begin{smallmatrix} + & 0\ -42 \end{smallmatrix}\big)$ \mathcal{L} *mq* (7)FSE(17)*^a* $|V_{\textit{\tiny{US}}}| = 0.2237\, (7)\, (^{+10}_{-\,0})_{m_{\!q}}(2)_{\textsf{FSE}}(4)_a$

- From $(A, \Phi, \Pi)_{\text{ini}}$ use molecular dynamics to move to $(A, \Phi, \Pi)_{\text{fin}}$
- Then do Monte Carlo accept/reject
- Rational Hybrid Monte Carlo of Clark and Kennedy approximates

- The RBC and UKOCD collaborations, using the numerically expensive dom t_{max} limit, too. For this reason lattice \mathcal{L} have two large simulations at different lattice spacings with physical light quarks. • The RBC and UKQCD collaborations, using the numerically expensive domain wall formulation, now it can be accounted for by simply replacing the momentum
- ^{*} Domain wall fermions preserve all the continuum symmetries of OCD a this technique has not been used in any work on light hadron *re* an are comp * Domain wall fermions preserve all the continuum symmetries of QCD at finite lattice spacing
- $*$ This is vital for measurements of many observables. * This is vital for measurements of many observables
- $\overline{1}$ $\overline{1}$ $\overline{0}$ $\overline{1}$ $\overline{0}$ $\overline{1}$ $I = \frac{11}{2}$ interacting field theory, the properties of a particle interaction in particle in partic a finite amplitude for a kaon to de **Spectroscopy that all hadron masses in a finite box 108** deviate from the set of α • For the first time, physical results have been produced for a process with 2 particles in the final state, probably \overline{p} above \overline{p} (\overline{p} \overline{p} and pseudoscalar methe amplitude for a kaon to decay to 2 pions (K $\rightarrow \pi\pi$) with the pions in an isospin 2 final state. Physical Review Letters, 108 (2012) 141601 and Physical Review D 86 (2012) 074513. $\mathcal{F}_{\mathcal{F}}$ and the proposition of view that the view that there is the view that there is the view that there is the view that the vi

Mira at ANL

Quantum Chromodynamics

• Relating standard model parameters to magazine and results concribedly required measured results generically requires knowing the value of a quark process inside a hadron (a matrix element) such

Computers, Algorithms and Software

• Physicists at Columbia built the 8,000 and 12,000 node QCDSP computers for QCD in 1997

• Columbia, RBRC and UK physicists built 3, 12,000 node QCDOC computers for QCD in 2004, working with IBM

T_A is the interval distribution ensembles with strange configuration ensembles with strange α and charm quark masses set at or very close to their physical Physics Results and Prospects were all cases turned out to be all cases turned out to be a good α

Major Development: Ensembles with Physical Quark Masses vsical Chiark Masses space of the the theory

- on the custom ASIC. These machines were built while IBM was producing the BG/L computer.
- Columbia (Christ, Kim) and Edinburgh (Boyle) physicists worked on the design of the BG/Q computer.
- Boyle has an extensively optimzed linear solver for the lattice Dirac equation that makes extensive use of the hardware features of the BG/Q. This solver was used in the chip design and testing stages to help validate the design.

Computers

Challenges and Prospects

The calculations reported here have been run on Mira at the ALCF, Sequoia at LLNL, the BGQ computers of the University of Edinburgh and the BGQ computers of BNL and the RBRC.

