Large-Scale Ab Initio Isobaric–Isothermal Simulations Using Self-Consistent van der Waals Inclusive Density Functionals Hsin-Yu Ko¹ Biswajit Santra¹ Robert A. DiStasio Jr.¹² Roberto Car¹

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Introduction

- Common Density Functional Theory (DFT) functionals do not include the non-local correlation effects that underlie van der Waals (vdW) or dispersion interactions, leading to serious errors for weakly bound systems.
- For systems with hydrogen bonds, the self-interaction error (SIE) of DFT is sizeable and hybrid DFT functionals can be utilized to alleviate this error.
- Here, we report a general-purpose self-consistent (SC) implementation of vdW-inclusive DFT functionals and discuss their effects on the equilibrium properties of several systems.

Nuclear quantum effects in pyridine I

Nuclear quantum effects (NQE) in pyridine I were investigated with **Path-Integral** (PI-AIMD) using the QE + i-PI implementation [6].

Nuclei	Volume	Volume	KE	EKE
	[Å ³ /molec]	Error	[eV/molec]	[eV/molec]
Classical	110.4	-0.99%	0.22	
Quantum	111.6	0.11%	1.19	0.97

Table 2: Equilibrium volumes and ionic kinetic energies (KE) from AIMD and PI-AIMD at the experimental thermodynamic conditions. Colored-noise generalized Langevin thermostats (with a Trotter dimension of 8) were used in all PI-AIMD simulations.

The calculated equilibrium volume including NQE is very close to

Effects of SC vdW on electronic structure

- The long-range vdW energy (E_{vdW}) typically represents a tiny fraction ($\sim 0.001\%$) of the total energy; hence, its influence on electronic properties are typically ignored.
- Utilizing SC vdW, we have found large effects on the electron density n(r) of molecular dimers, alkali-metal dimers, transition-metal surfaces, and organic-metal interfaces [1].
- Methodology: modified Kohn-Sham effective potential [2].

 $v[n(r)] = v_{GGA/hybrid}[n(r)] + \frac{\delta E_{vdW_{SC}^{TS}}[n(r)]}{\delta n(r)}$ $E_{vdW_{\rm SC}^{\rm TS}}[n(r)] = -\sum_{A < B} f_{damp}^{AB}[n(r)] \frac{C_{6AB}[n(r)]}{R_{AB}^6}$

vdW inclusive functionals are implemented in FHI-AIMS [3] and 3 Quantum ESPRESSO (QE) [4].

N	letal	PBE	PBE+vdW _{SC} ^{surf}	Expt.
	Cu	4.89	4.95	4.94
	Rh	5.26	5.55	5.60
	hΑ	4 4 4	4 74	4 74



the AIMD volume at p_{exp} and T_{exp} , in spite of the large excess ionic kinetic energy (EKE).

- NQE are sizeable in the intramolecular motions but small in the intermolecular motions; equilibrium volumes are primarily determined by intermolecular motions.
- To lowest-order in \hbar , NQE in the intermolecular motions are equivalent to the addition of a ΔT correction to the simulation temperature **T**,

$$\Delta T = \frac{\hbar^2}{36 M T^2} \left\langle \left| \vec{F} \right|^2 \right\rangle.$$

 $T \Delta T$ Pyridine I 153 6.8 Pyridine II 293 5.5 Table 3: Quantum temperature corrections $(\Delta T \text{ in Kelvin}) \text{ using (1).}$

Here, *M* is the mass of one pyridine molecule and \vec{F} is the force on its center

of mass.

The small magnitudes of ΔT reveal the quasi-classical character of the intermolecular motions at T_{exp} .

(1)

Ice triple point with PBE0+vd W_{SC}^{TS}



 $\mathsf{PBE0} + vdW_{SO}^{T}$

Table 1: Work functions (in eV) of various metal (111) surfaces obtained from experiment and theory.

Figure 1: Integrated electron density differences $[\Delta n(r)]$ for the methane (CH₄) dimer separated by a C-C distance of 6.72 Å along the *z* axis.

Role of vdW and anharmonicity on the equilibrium structures of molecular crystals



PBE+vdW $_{sc}^{TS}$





Figure 4: Snapshots of ice I*h*, II, and III at the experimental triple point (0.21 GPa and 238 K).

- ► These large-scale *NpT* simulations were made possible by the linear-scaling exact $\sum_{n=40}^{\infty}$ exchange (EXX) and vdW algorithms im-田 -60 plemented in QE [7]. The implementation uses analytical EXX stress tensors.
- Triple point properties calculated proved systematically from PBE PBE+vdW^{TS}_{SC} and to PBE0+vdW^{TS}_{SC}.



Figure 5: Volume (ΔV) IMand enthalpy (ΔH) to differences between the ice phases at the triple point.

Conclusions

In this work, we have provided a general-purpose implementation of vdW-inclusive DFT functionals for variable-cell NpT AIMD simulations and report their successful application to a wide variety of systems of interest in biology, chemistry, and physics.

mizes the dipolar interaction en-

P412121 P21/n Pnma P2₁/n P2₁/n P2₁2₁2₁

Figure 3: (a) PLMC structures from ergy. AIMD at p_{exp} and both T = 0 K and • NpT dynamics captures anhar- $T = T_{exp}$. (b) PLMC structures from monicity and the SC vdW con-AIMD at p_{exp} and T_{exp} with and without vdW. The inset shows an overlay of the tribution to the stress tensor has calculated equilibrium structure with the been implemented in QE. experimental X-ray structure.

• The anharmonicity is large (with the PBE+vdW $_{SC}^{TS}$ pyridine volume at $T = 153 \text{ K} \sim 4\%$ larger than at T = 0 K). Classical simulations seem to be sufficient to capture the anharmonicity in this system.

References and acknowledgements

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