

Wick Contractions for Nucleon-Nucleon Scattering and Matrix Elements

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Introduction

Automatic Contraction Generation

- A Wick contraction is a way to tie creation and annihilation operators into propagators. Each contraction corresponds to a Feynman diagram.
- The proton-proton scattering by the $\Delta I = 2$ weak operator insertion has 4896 contractions. These reduce to 2208 contractions for degenerate up- and down-quark masses.
- Each contraction comes with a sign. These signs are
 - ▶ important for the physics: the signs control the interference between the diagrams. tricky to keep track of by hand: they sensitively depend on operator ordering.
 - absolutely critical to get right: a single error casts the whole result into doubt.

Baryon Block Formalism

Once we have done the Wick contractions, we are left with a number of

It is impractical to write code for 2208 contractions by hand.

- If calculated by hand, the chance of 2208 correct signs would be low.
- We need software to help write the correct code.
- A good contraction generator would make this tedious task mindless. ▶ It should:
 - Parse familiar physics expressions,
 - ▶ Be easy to extend,
 - Algorithmically generate all Wick contractions,
 - Generate code for a variety of lattice QCD software.

Generation in Mathematica

Mathematica makes it easy to parse and manipulate

quark propagators whose indices need to be tied up appropriately. We can take advantage of some physics knowledge to simplify and speed up the contractions. For example, we know that baryons are color singlets. This lets us build a baryon block,

 $BB^{\mu}(x_{f}, \{x_{i1}, x_{i2}, x_{i3}\})^{\beta\delta\phi}_{bdf} = \sum \epsilon^{ace} \Gamma_{\alpha\gamma} S(x_{f}, x_{i1})^{\alpha\beta}_{ab} S(x_{f}, x_{i2})^{\gamma\delta}_{cd} S(x_{f}, x_{i2})^{\mu\phi}_{ef}$

which is color neutral at the sink.

- We can factor every contraction into baryon blocks, possibly some left-over quark propagators, and constant, sparse tensors (which we can pre-compute). See below for a simple example.
- All nucleon-nucleon partial-wave scattering can ultimately be reduced to a handful of sums like

$$\sum_{i,j} BB^i(x + \Delta x) BB^j(x) T_{ij}$$

where *i*, *j* are collective indices, and T is the tensor which tells us how to connect them.

Sources

In order to measure scattering in particular partial waves, we need

- symbolic expressions, and is interactive and fast enough to allow error correction and fast iteration.
- We have written an automatic contraction generator using Mathematica.



- Proton two-point function has two Wick contractions \nearrow .
 - Begin with physics expression: $\langle p^{\mu}(x_f) | p_{\mu}(x_i) \rangle$
 - Translate into simple mathematica expression: $pp = Proton[sink][xf, \mu] ** bar[Proton[source][xi, \nu]] ** SpinProjector[Same, Spin[\nu, \mu]];$
 - Symbolically generate contractions: contractions = Contract[pp]; Notation [contractions]
 - $\operatorname{Same}_{\nu\mu} \operatorname{Cy}_{5_{\nu'}} \operatorname{Cy}_{5_{\nu'}} \operatorname{Cy}_{5_{\nu'}} \operatorname{S}[\operatorname{down}, \operatorname{xf} - \operatorname{xi}]_{\nu'\nu'}^{\mathfrak{b}'\mathfrak{b}'} \operatorname{S}[\operatorname{up}, \operatorname{xf} - \operatorname{xi}]_{\mu\nu}^{\mathfrak{a}'\mathfrak{a}'} \operatorname{S}[\operatorname{up}, \operatorname{xf} - \operatorname{xi}]_{\rho'\rho'}^{\mathfrak{c}'\mathfrak{c}'} e^{\mathfrak{a}'\mathfrak{b}'\mathfrak{c}'} e^{\mathfrak{a}'\mathfrak{b}'\mathfrak{b}'} e^{\mathfrak{a}'\mathfrak{b}'\mathfrak{b}'} e^{\mathfrak{a}'\mathfrak{b}'\mathfrak{b}'} e^{\mathfrak{a}'\mathfrak{b}'\mathfrak{b}'} e^{\mathfrak{a}'\mathfrak{b}'\mathfrak{b}'} e^{\mathfrak{a}'\mathfrak{b}'\mathfrak{b}'} e^{\mathfrak{a}'\mathfrak{b}'\mathfrak{b}'} e^{\mathfrak{a}'\mathfrak{b}'\mathfrak{b}'} e^{\mathfrak{a}'\mathfrak{b}'\mathfrak{b}'} e^{\mathfrak{a}'\mathfrak{b}'} e^{\mathfrak{a}'} e^{\mathfrak{a}'} e^{\mathfrak{a}'} e^{\mathfrak{a}'} e^{\mathfrak{a}'\mathfrak{b}'} e^{\mathfrak{a}'} e^{\mathfrak{$
 - Take generated contractions and generate code
 - using QDP primitives Generate [QDP] [contractions]

-colorTrace(spinTrace(S_up_xf_from_xi * projectorSame) * spinTrace(quarkContract24(S_up_xf_from_xi, S_down_xf_from_xi * gammaC5) * gammaC5)) -trace(gammaC5 * S_up_xf_from_xi * projectorSame * quarkContract24(S_up_xf_from_xi, S_down_xf_from_xi * gammaC5))

for creating tensors for the baryon block formalism Generate [BaryonBlocks] [contractions] // Notation

for_color(b) { for_color(d) { for_color(f) { for_spin(epsilon) { for_spin(beta) {

- sources with good overlap with the corresponding initial spatial wavefunction.
- ► For *p*-wave (and higher-order) scattering, this requires sources where the baryons are spatially separated.
- We put each baryon on a single lattice site, but displace the sites relative to one another.

 $\overline{N}(x + \Delta x, t)\overline{N}(x, t)$

- $\Delta x = 0$ has good overlap with the *s*-wave (and some to L = 4 and higher) but none with the *p*-wave
- $\Delta x \propto \pm \hat{x}, \pm \hat{y}, \pm \hat{z}$ allows us to construct in post-processing
- nonlocal sources with good s-wave overlap (via the isotropic combination)
- nonlocal sources with good p-wave overlap (via antisymmetric combinations)

Sinks

- Similarly, we need sinks with good overlap with the corresponding final wavefunction.
- For partial-wave scattering, we store the full momentum dependence of each correlator. This makes it cheap to try any desired sink, because we can pull out particular displacements

for_spin(delta) { phi) { for_spin = 0; temp -= epsilon3(b, d, f) * diquark[beta][phi] * spin_projector_Same[delta][epsilon]; -= epsilon3(b, d, f) * diquark[beta][delta] * spin_projector_Same[phi][epsilon]; tensor_singleProton.add(temp, tensor_index(b, d, f, epsilon, beta, delta, phi)); } } } } $-\left(\mathsf{Same}_{\phi \in} \mathsf{C}\gamma_{5_{\beta} \delta} + \mathsf{Same}_{\delta \in} \mathsf{C}\gamma_{5_{\beta} \phi}\right) \mathsf{BB}^{\epsilon}\left[\mathsf{xf,} \{\mathsf{xi, xi, xi}\}\right]_{\mathsf{bdf}}^{\beta \delta \phi} \epsilon^{\mathsf{bdf}}$

Using CUDA

The two-baryon contraction problem can be viewed as a calculation of the dense-sparse-dense bilinear product

 $BB(x + \Delta x)^T \cdot T \cdot BB(x)$

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- using indirect addressing, we have to compute for all sites x a product such as v[I1[k]] * T[k] * w[I2[k]] and sum over k.
- ▶ compute 11, 12 once and store them along with v, w and T in global GPU memory
- ▶ read v[I1[k]], w[I2[k]] and T[k] from global memory and compute the product, one k per



1 7 -1 6 -2 8 5 17 -3 9 7 13 11 2 2

6 -2 8 5 17 -3 9 7 13 11 2 2



or fold in any spatial wavefunction.

For nucleon-nucleon matrix elements, this approach is prohibitively expensive, and so we fix our sinks in the same way we fix our sources.

Operator Insertion

- The parity-violating process we are interested in conserves momentum, and thus our operator must be projected to zero momentum. This can be accomplished
- Stochastically (allowing us to use arbitrary sinks), at the cost of increased noise. Exactly (which requires a sum over the lattice volume). This is feasible with fixed sinks.

thread

perform efficient binary reduction ▶ for each x, create separate CUDA stream to maximize concurrency host thread concurrently computes all FFTs and baryon blocks relevant for next step of the calculation

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