

# **Water, soft matter and reactions in solution: major challenges to microscopic modeling and simulation**

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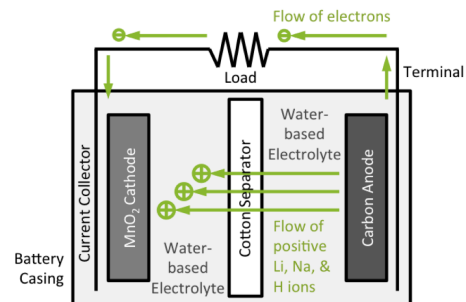
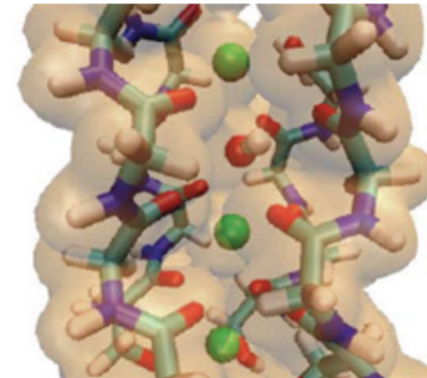
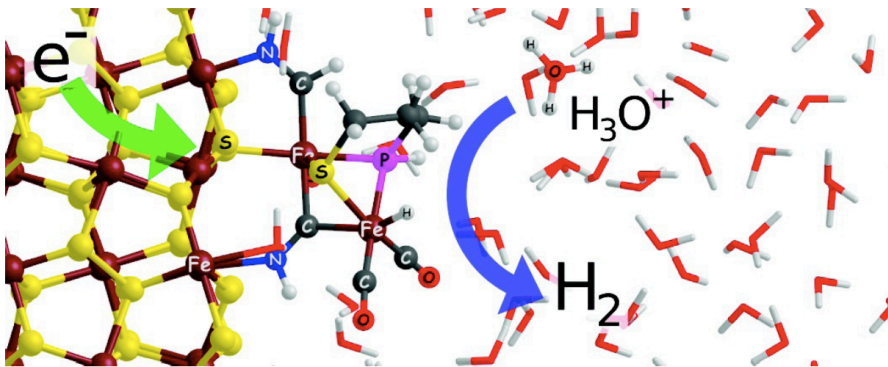
*2013 SciDAC-3 Principal Investigator Meeting, July 24-26, 2013*

# Advanced Modeling of Ions in Solutions, on Surfaces, and in Biological Environments



DOE SciDAC

<http://amis-scidac.org/>





## The **AMIS** team

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# **Electronic structure, ab-initio molecular dynamics, spectral properties**

- **The Standard Model**

Density functional theory for ground–state electronic structure

GW-BSE for electronic excitation and spectral properties (photoemission, optical, uv, x-ray absorption... )

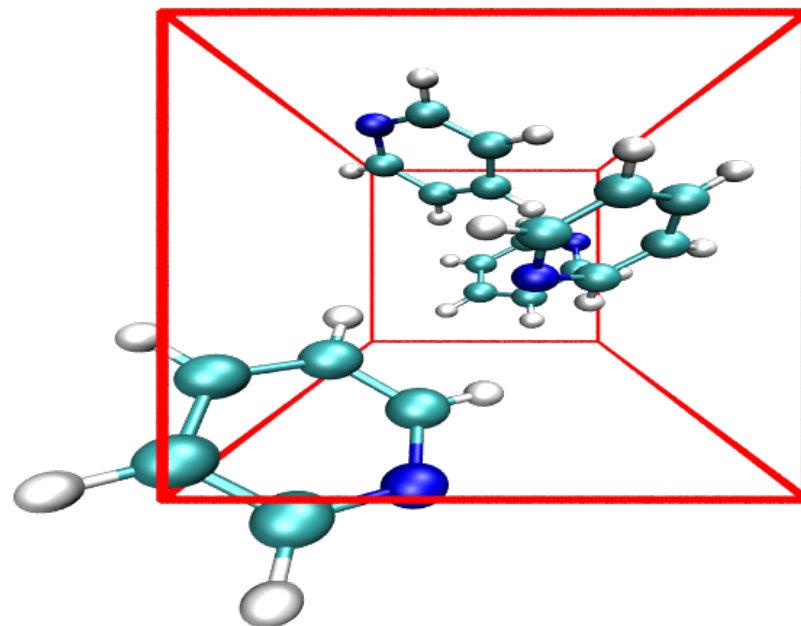
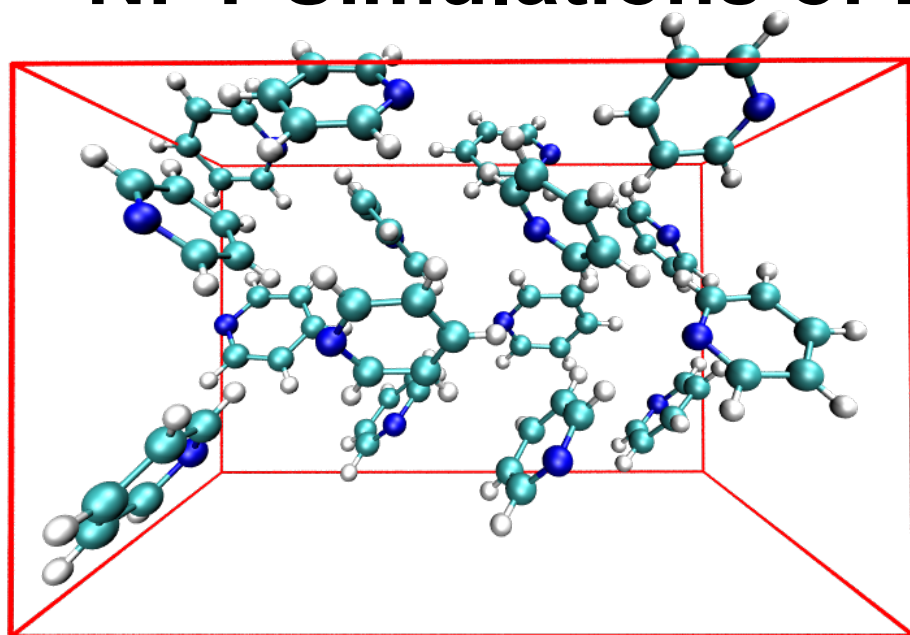
- **Beyond the Standard Model**

# Weak interactions (H-bonding, vdW)

DFT functionals must be enhanced to model vdW interactions (an electron correlation effect)

H-bonding is described qualitatively correctly by common functionals but issues of accuracy remain (self-interaction error & overestimation of H-bond strength)

# NPT Simulations of Pyridine Polymorphs



	<b>h5-I (1 Bar)</b> Pna2 <sub>1</sub> (Orthorhombic)				<b>h5-II (1.1 GPa)</b> P2 <sub>1</sub> 2 <sub>1</sub> 2 <sub>1</sub> (Orthorhombic)			
	0 K		153 K		0 K		293 K	
	Volume [Å <sup>3</sup> ]	Error	Volume [Å <sup>3</sup> ]	Error	Volume [Å <sup>3</sup> ]	Error	Volume [Å <sup>3</sup> ]	Error
<b>Exp</b>	--		1784		--		414	
<b>PBE</b>	2185	+22.5%	2250 ± 24	+26.1 %	442	+6.8 %	493 ± 8	+19.1 %
<b>PBE+vdw</b>	1710	- 4.2 %	1770 ± 22	- 0.7 %	396	- 4.3 %	414 ± 7	+0.01 %

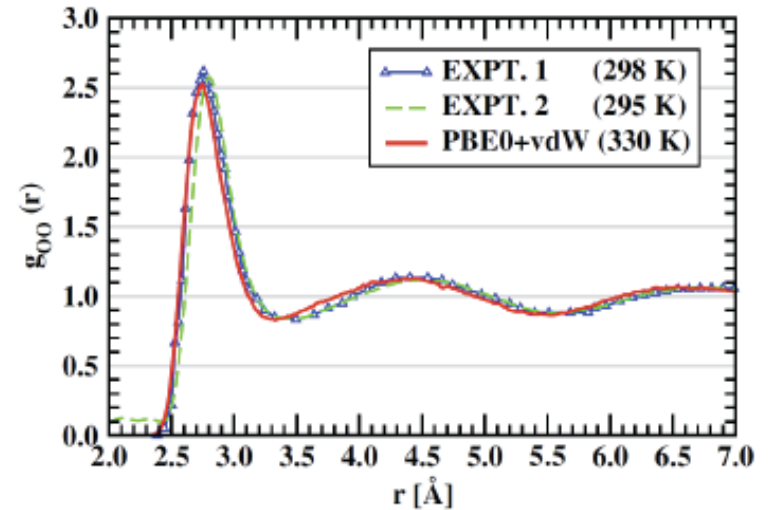
vdW functional: Tkatchenko, Scheffler, *PRL* (2009); Tkachenko, DiStasio, Car, Scheffler, *PRL* (2012)

From H-Y Ko, R. DiStasio, RC (to be published)

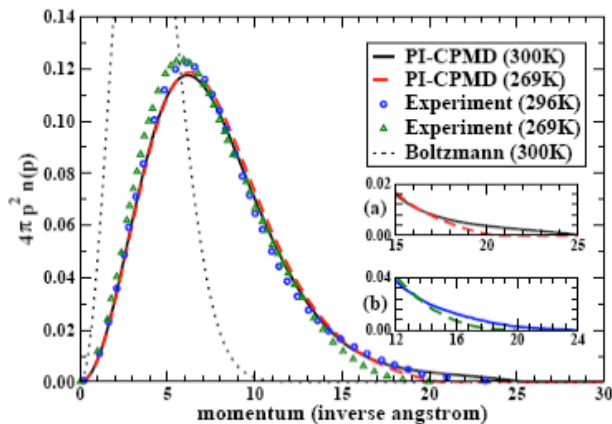
# Water: DFT approximation, quantum nuclei

$(\text{H}_2\text{O})_N$	PBE (330K)	PBE+vdW (330K)	Expt.
32	0.78 (40 ps)	0.98 (50 ps)	-
64	0.80 (40 ps)	0.98 (30 ps)	-
128	0.79 (30 ps)	0.99 (20 ps)	0.99
ice	0.97	1.01	0.93

Table : Densities (in  $\text{g}/\text{cm}^3$ ) resulting from AIMD-NPT simulations of liquid water at the PBE and PBE+vdW levels of theory. All simulations were run using  $E_0 = 130Ry$  ( $E_{cut} = 150Ry$ ). The simulation temperature was set to  $T = 330\text{K}$  to approximately account for NQE.



From B. Santra, R. DiStasio, Z. Li, R.C. (to be published)

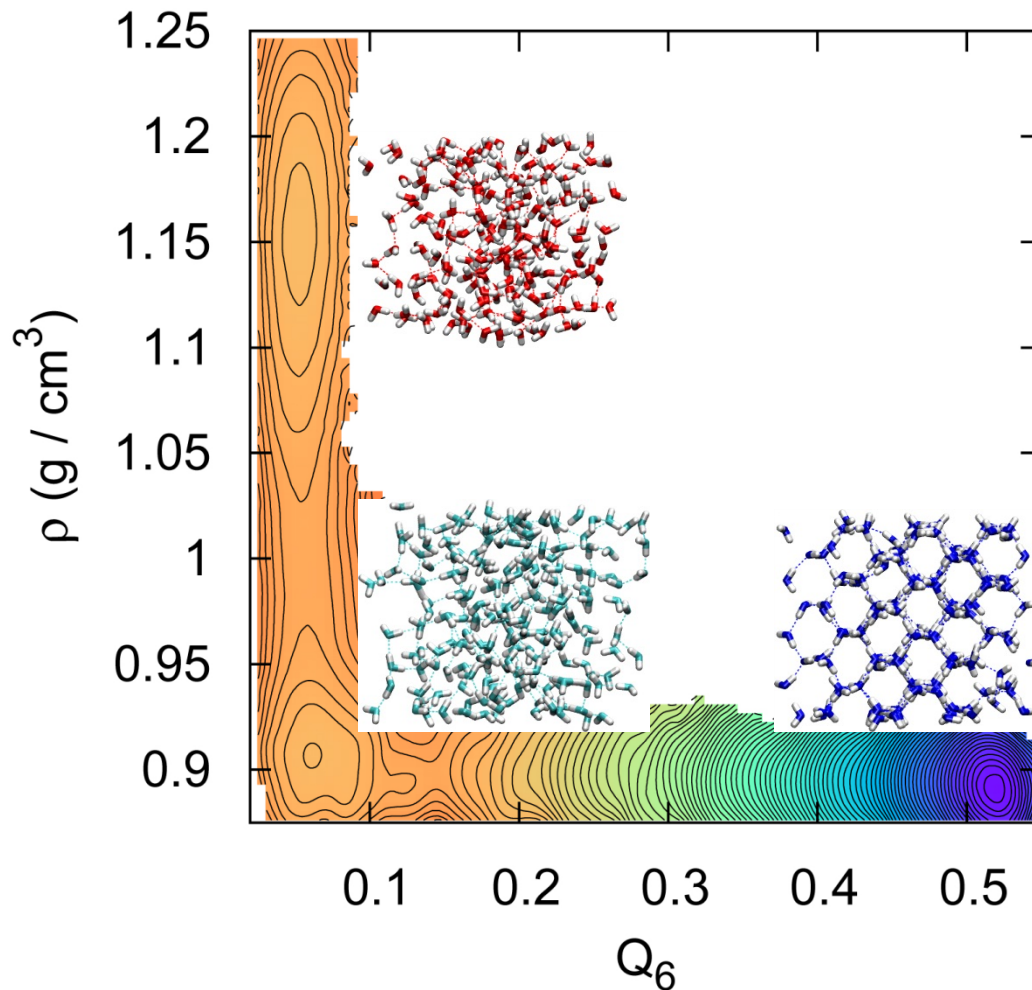


J.Morrone, R.C., *PRL* (2008)

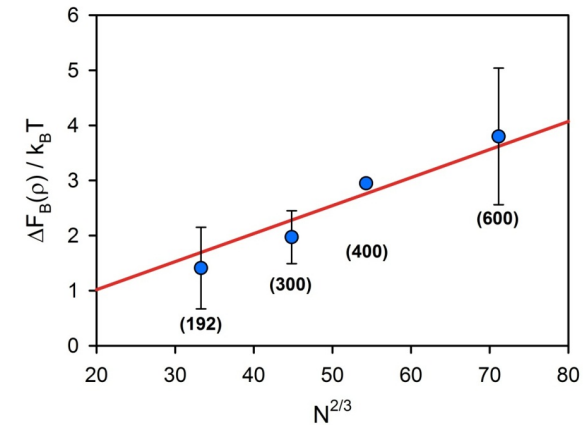
Soon we will start extensive PI-AIMD simulations on BG-Q (ALCC grant) using PBE0+vdW and color noise thermostats (Ceriotti, Parrinello); very effective parallelization

# Anomalies, Phase diagram, Sampling Issues

- Secondary critical point (HDL-LDL transition)
- Is the transition genuinely first-order?



Finite size scaling with well tempered metadynamics



J. Palmer, R. Car, P. Debenedetti,  
Faraday Discuss, (2013)



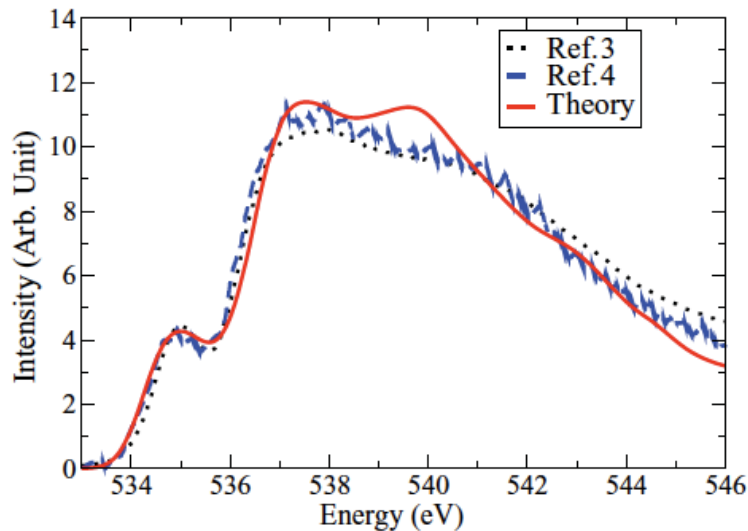
# **The large scale challenge**

Collaboration with FASTMATH on Pole Expansion  
and Selected Inversion Method for Accelerating  
Electronic Structure Calculations

# Electron excitations

GW-BSE + good sampling of disordered structures

Ionization potential of water ions (poster by C.W. Swartz and X. Wu, Temple)



x-ray absorption spectra  
of liquid water (GW:  
static COHSEX)  
(L. Kong, X. Wu, R.C.  
*PRB* (2012))

# DFT Issues

- Functional fix not always works, plethora of functionals, no systematic procedure for improving
- Self-interaction error may cause severe difficulties, for example in the case of the water radical OH and in general when dealing with chemical reactions
- Incorrect dissociation of molecules into non-interacting fragments (“static” or strong correlation)

# Occupation Probabilities-Natural Orbital Functional Theory

$$E = \langle \Psi | H | \Psi \rangle = \text{Tr} [ H \pi ]$$

$$\Psi(x_1, x_2, \dots, x_N) = \sum_{\mathbf{n}} C_{\mathbf{n}} \Phi_{\mathbf{n}}(x_1, x_2, \dots, x_N) \quad (\mathbf{n} \equiv n_1 n_2 \dots n_{N/2})$$

Simplest version: Doubly Occupied NO states – seniority 0 (DOCI)

$$\pi(x'_1, x'_2; x_1, x_2) = N(N-1) \int dx_3 \cdots dx_N \Psi(x'_1, x'_2, \dots, x_N) \Psi^*(x_1, x_2, \dots, x_N),$$

OP-NOFT: R Gebauer, MH Cohen, R.C.

$$E = 2 \sum_i h_{ii} p_1(i) + \sum_{ij} p_{11}(ij) [2J_{ij} - K_{ij}] + \sum_{i \neq j} p_{10}^{1/2}(ij) p_{01}^{1/2}(ij) s(i) s(j) \xi(ij) K_{ij}$$

$$J_{ij} \equiv \langle \varphi_i \varphi_j | v | \varphi_i \varphi_j \rangle = \int d\mathbf{r}_1 d\mathbf{r}_2 \frac{\varphi_i(\mathbf{r}_1) \varphi_j(\mathbf{r}_2) \varphi_i(\mathbf{r}_1) \varphi_j(\mathbf{r}_2)}{r_{12}} \quad 0 \leq \xi(ij) \leq 1$$

$$K_{ij} \equiv \langle \varphi_i \varphi_j | v | \varphi_j \varphi_i \rangle = \int d\mathbf{r}_1 d\mathbf{r}_2 \frac{\varphi_i(\mathbf{r}_1) \varphi_j(\mathbf{r}_2) \varphi_j(\mathbf{r}_1) \varphi_i(\mathbf{r}_2)}{r_{12}}$$

$$p_1(k) \equiv \sum_{\mathbf{n}} C_{\mathbf{n}}^2 \nu_{k,\mathbf{n}} \quad p_{11}(ij) \equiv \sum_{\mathbf{n}} C_{\mathbf{n}}^2 \nu_{i,\mathbf{n}} \nu_{j,\mathbf{n}} \quad p_{10}(ij) \equiv \sum_{\mathbf{n}} C_{\mathbf{n}}^2 \nu_{\mathbf{n},i} (1 - \nu_{\mathbf{n},j})$$

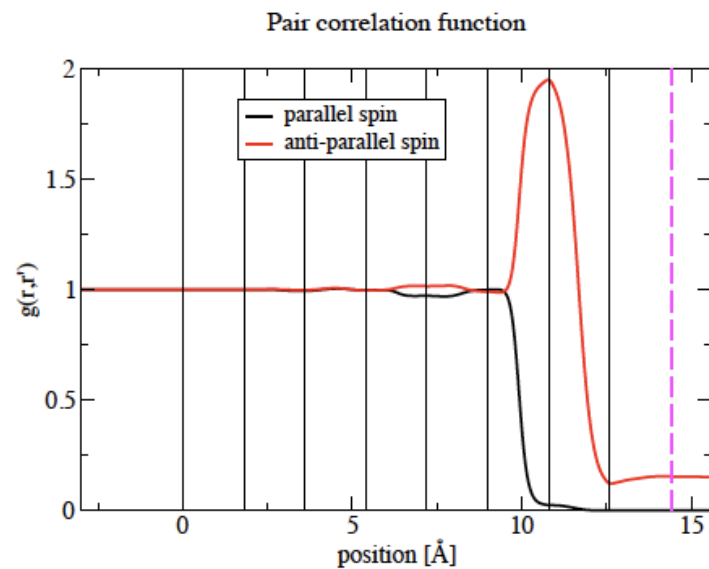
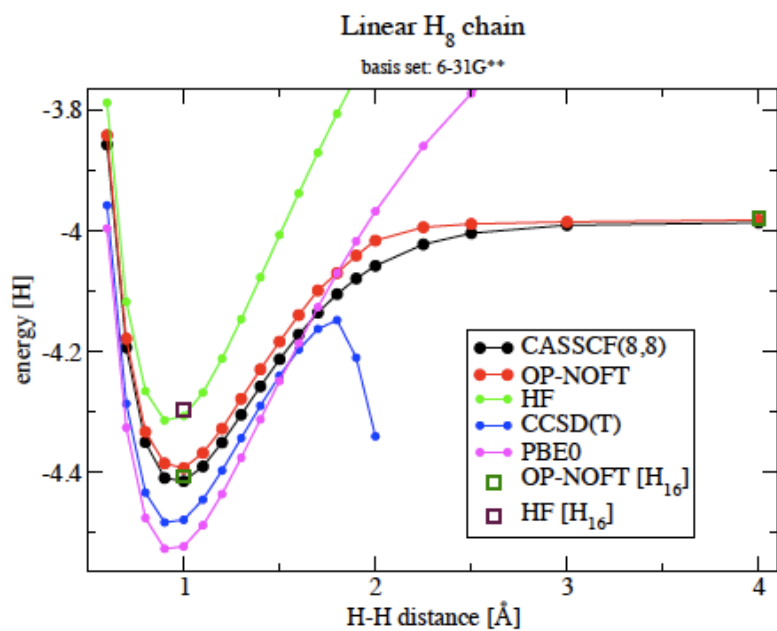
$$p_1(i) = p_{11}(ij) + p_{10}(ij)$$

$$\xi(ij) \approx F[p_{11}] \quad \text{while exact } \xi(ij) = F[p_{11\dots 1}]$$

$p_1, p_{11}$  satisfy bounds and sum rules to ensure  $N$ -representability

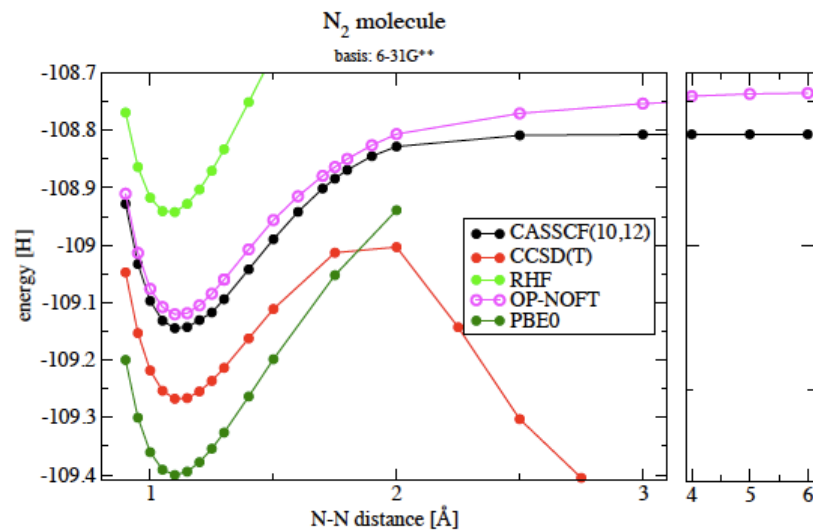
The ground-state energy is the infimum of  $E$  wrt  $\varphi_i(\mathbf{r}), p_1(i), p_{11}(ij)$  subject to constraints, bounds and sum rules

Minimization of  $E$  has the same scaling with  $N$  of Hartree-Fock theory

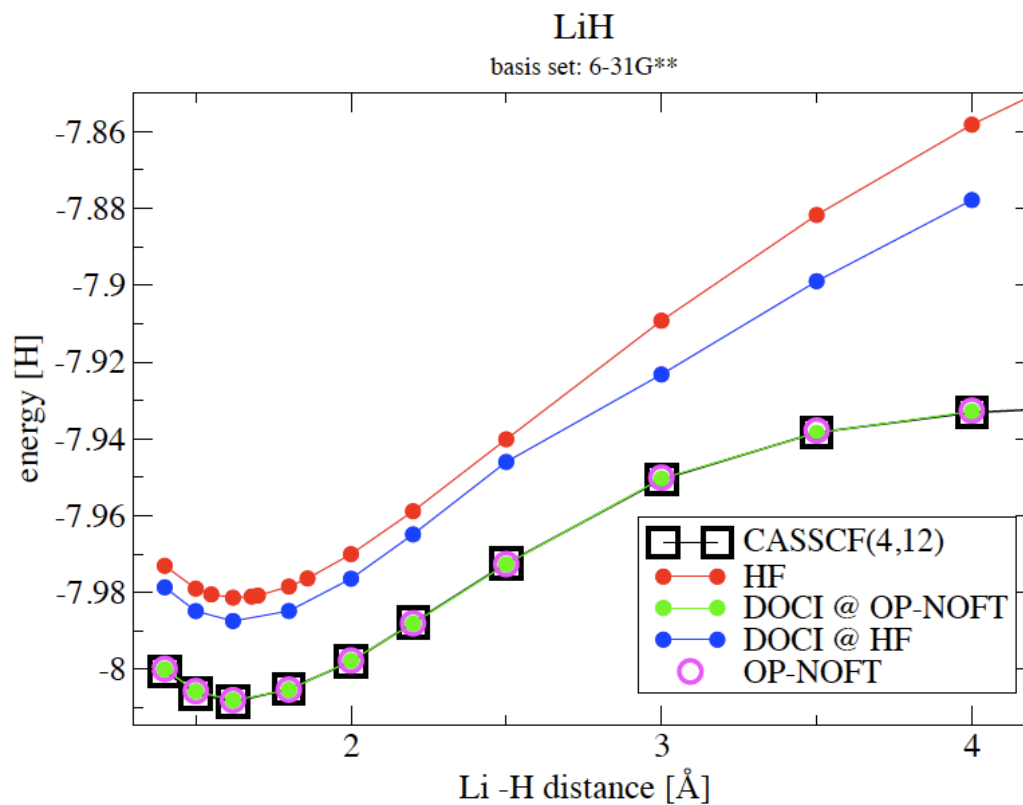


$$\int dr_2 \pi(r', r_2; r, r_2) = (N - 1)\rho(r, r') \text{ and}$$

$$\int dr_1 dr_2 \pi(r_1, r_2; r_1, r_2)w(r_{12}) \geq 0,$$



# OP-NOFT vs DOCI



DOCI: Weinhold and Bright Wilson (1967)

- Calculation of forces, structural optimization and AIMD “straightforward”
- This version of the theory (corresponding to DOCI) includes only *intra-shell* correlations among electrons of opposite spins
- A more accurate version including also *inter-shell* correlations among electrons of arbitrary spin is possible: the scaling would still be polynomial but worse than HF



## **Concluding Remarks**

**SIZE, TIME (SAMPLING), ACCURACY** are all very important issues in molecular simulations

**The goal of chemical accuracy in condensed phase simulations can be within reach, by combining physics theory, effective algorithms and high performance computer platforms**